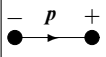


7.4 Fields associated with media

Polarisation

Definition of electric dipole moment	$\mathbf{p} = q\mathbf{a}$	(7.80)	$\pm q$ end charges $\mathbf{a}$ charge separation vector (from $-$ to $+$)	
Generalised electric dipole moment	$\mathbf{p} = \int_{\text{volume}} \mathbf{r}' \rho \, d\tau'$	(7.81)	$\mathbf{p}$ dipole moment ρ charge density $d\tau'$ volume element $\mathbf{r}'$ vector to $d\tau'$	
Electric dipole potential	$\phi(\mathbf{r}) = \frac{\mathbf{p} \cdot \mathbf{r}}{4\pi\epsilon_0 r^3}$	(7.82)	ϕ dipole potential $\mathbf{r}$ vector from dipole ϵ_0 permittivity of free space	
Dipole moment per unit volume (polarisation) ^a	$\mathbf{P} = n\mathbf{p}$	(7.83)	$\mathbf{P}$ polarisation n number of dipoles per unit volume	
Induced volume charge density	$\nabla \cdot \mathbf{P} = -\rho_{\text{ind}}$	(7.84)	ρ_{ind} volume charge density	
Induced surface charge density	$\sigma_{\text{ind}} = \mathbf{P} \cdot \hat{\mathbf{s}}$	(7.85)	σ_{ind} surface charge density $\hat{\mathbf{s}}$ unit normal to surface	
Definition of electric displacement	$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$	(7.86)	$\mathbf{D}$ electric displacement $\mathbf{E}$ electric field	
Definition of electric susceptibility	$\mathbf{P} = \epsilon_0 \chi_E \mathbf{E}$	(7.87)	χ_E electrical susceptibility (may be a tensor)	
Definition of relative permittivity ^b	$\epsilon_r = 1 + \chi_E$ $\mathbf{D} = \epsilon_0 \epsilon_r \mathbf{E}$ $= \epsilon \mathbf{E}$	(7.88) (7.89) (7.90)	ϵ_r relative permittivity ϵ permittivity	
Atomic polarisability ^c	$\mathbf{p} = \alpha \mathbf{E}_{\text{loc}}$	(7.91)	α polarisability $\mathbf{E}_{\text{loc}}$ local electric field	
Depolarising fields	$\mathbf{E}_d = -\frac{N_d \mathbf{P}}{\epsilon_0}$	(7.92)	$\mathbf{E}_d$ depolarising field N_d depolarising factor =1/3 (sphere) =1 (thin slab $\perp$ to $\mathbf{P}$) =0 (thin slab $\parallel$ to $\mathbf{P}$) =1/2 (long circular cylinder, axis $\perp$ to $\mathbf{P}$)	
Clausius–Mossotti equation ^d	$\frac{n\alpha}{3\epsilon_0} = \frac{\epsilon_r - 1}{\epsilon_r + 2}$	(7.93)		

^a Assuming dipoles are parallel. The equivalent of Equation (7.112) holds for a hot gas of electric dipoles.

^b Relative permittivity as defined here is for a linear isotropic medium.

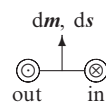
^c The polarisability of a conducting sphere radius a is $\alpha = 4\pi\epsilon_0 a^3$. The definition $\mathbf{p} = \alpha\epsilon_0 \mathbf{E}_{\text{loc}}$ is also used.

^d With the substitution $\eta^2 = \epsilon_r$ [cf. Equation (7.195) with $\mu_r = 1$] this is also known as the “Lorentz–Lorenz formula.”



Magnetisation

Definition of magnetic dipole moment	$d\mathbf{m} = I d\mathbf{s}$	(7.94)	$d\mathbf{m}$ dipole moment I loop current $d\mathbf{s}$ loop area (right-hand sense with respect to loop current)
Generalised magnetic dipole moment	$\mathbf{m} = \frac{1}{2} \int_{\text{volume}} \mathbf{r}' \times \mathbf{J} d\tau'$	(7.95)	$\mathbf{m}$ dipole moment $\mathbf{J}$ current density $d\tau'$ volume element $\mathbf{r}'$ vector to $d\tau'$
Magnetic dipole (scalar) potential	$\phi_m(\mathbf{r}) = \frac{\mu_0 \mathbf{m} \cdot \mathbf{r}}{4\pi r^3}$	(7.96)	ϕ_m magnetic scalar potential $\mathbf{r}$ vector from dipole μ_0 permeability of free space
Dipole moment per unit volume (magnetisation) ^a	$\mathbf{M} = n\mathbf{m}$	(7.97)	$\mathbf{M}$ magnetisation n number of dipoles per unit volume
Induced volume current density	$\mathbf{J}_{\text{ind}} = \nabla \times \mathbf{M}$	(7.98)	$\mathbf{J}_{\text{ind}}$ volume current density (i.e., A m ⁻²)
Induced surface current density	$\mathbf{j}_{\text{ind}} = \mathbf{M} \times \hat{\mathbf{s}}$	(7.99)	$\mathbf{j}_{\text{ind}}$ surface current density (i.e., A m ⁻¹) $\hat{\mathbf{s}}$ unit normal to surface
Definition of magnetic field strength, $\mathbf{H}$	$\mathbf{B} = \mu_0(\mathbf{H} + \mathbf{M})$	(7.100)	$\mathbf{B}$ magnetic flux density $\mathbf{H}$ magnetic field strength
Definition of magnetic susceptibility	$\mathbf{M} = \chi_H \mathbf{H}$	(7.101)	χ_H magnetic susceptibility. χ_B is also used (both may be tensors)
	$= \frac{\chi_B \mathbf{B}}{\mu_0}$	(7.102)	
	$\chi_B = \frac{\chi_H}{1 + \chi_H}$	(7.103)	
Definition of relative permeability ^b	$\mathbf{B} = \mu_0 \mu_r \mathbf{H}$	(7.104)	μ_r relative permeability μ permeability
	$= \mu \mathbf{H}$	(7.105)	
	$\mu_r = 1 + \chi_H$	(7.106)	
	$= \frac{1}{1 - \chi_B}$	(7.107)	



^aAssuming all the dipoles are parallel. See Equation (7.112) for a classical paramagnetic gas and page 101 for the quantum generalisation.

^bRelative permeability as defined here is for a linear isotropic medium.



Paramagnetism and diamagnetism

Diamagnetic moment of an atom	$\mathbf{m} = -\frac{e^2}{6m_e} Z \langle r^2 \rangle \mathbf{B} \quad (7.108)$	$\mathbf{m}$ magnetic moment $\langle r^2 \rangle$ mean squared orbital radius (of all electrons) Z atomic number $\mathbf{B}$ magnetic flux density m_e electron mass $-e$ electronic charge
Intrinsic electron magnetic moment ^a	$\mathbf{m} \simeq -\frac{e}{2m_e} g \mathbf{J} \quad (7.109)$	$\mathbf{J}$ total angular momentum g Landé g-factor (=2 for spin, =1 for orbital momentum)
Langevin function	$\mathcal{L}(x) = \coth x - \frac{1}{x} \quad (7.110)$ $\simeq x/3 \quad (x \lesssim 1) \quad (7.111)$	$\mathcal{L}(x)$ Langevin function
Classical gas paramagnetism ($ \mathbf{J} \gg \hbar$)	$\langle M \rangle = nm_0 \mathcal{L} \left(\frac{m_0 B}{kT} \right) \quad (7.112)$	$\langle M \rangle$ apparent magnetisation m_0 magnitude of magnetic dipole moment n dipole number density T temperature k Boltzmann constant
Curie's law	$\chi_H = \frac{\mu_0 n m_0^2}{3kT} \quad (7.113)$	χ_H magnetic susceptibility
Curie-Weiss law	$\chi_H = \frac{\mu_0 n m_0^2}{3k(T - T_c)} \quad (7.114)$	μ_0 permeability of free space T_c Curie temperature

^aSee also page 100.

Boundary conditions for E , D , B , and H ^a

Parallel component of the electric field	$E_{\parallel}$ continuous	(7.115)	$\parallel$ component parallel to interface
Perpendicular component of the magnetic flux density	$B_{\perp}$ continuous	(7.116)	$\perp$ component perpendicular to interface
Electric displacement ^b	$\hat{s} \cdot (\mathbf{D}_2 - \mathbf{D}_1) = \sigma$	(7.117)	$\mathbf{D}_{1,2}$ electrical displacements in media 1 & 2 $\hat{s}$ unit normal to surface, directed 1 $\rightarrow$ 2 σ surface density of free charge
Magnetic field strength ^c	$\hat{s} \times (\mathbf{H}_2 - \mathbf{H}_1) = \mathbf{j}_s$	(7.118)	$\mathbf{H}_{1,2}$ magnetic field strengths in media 1 & 2 $\mathbf{j}_s$ surface current per unit width



^aAt the plane surface between two uniform media.

^bIf $\sigma = 0$, then $D_{\perp}$ is continuous.

^cIf $\mathbf{j}_s = \mathbf{0}$ then $H_{\parallel}$ is continuous.

